

Abstracts for the talks

Invited talk abstracts

1. Daniel Lidar (USC, USA; remote talk)

Title: Scaling Advantage in Approximate Optimization with Quantum Annealing

Abstract:

This talk will begin with a review of a decade of efforts to demonstrate a scaling advantage in optimization using quantum annealing. Exact optimization has proven to be an elusive target, but recent work has finally demonstrated a quantum scaling advantage in approximate optimization. Tailored quantum error suppression and correction methods play an important role in this demonstration. The advantage is achieved for a certain class of spin-glass problems, where, for sufficiently large optimality gaps, quantum annealing demonstrates a time-to-approximate solution that scales better than PT-ICM, the state-of-theart classical method.

This is joint work with Humberto Munoz-Bauza, arXiv:2401.07184

2. Alexander Whiticar (D-Wave Systems Inc.)

Coherence advancements for next-generation quantum technologies

D-Wave Quantum continually works on developing next-generation superconducting quantum computing technologies. Recent results from a new quantum annealing prototype have brought into focus the importance of realizing and augmenting the role of coherent quantum dynamics during operation [1,2]. To that end, D-Wave Quantum has maintained an intensive in-house materials science program to reduce sources of noise in our quantum processors. In this talk I will discuss a broad range of methods to measure flux and charge noise

in flux-like qubits. It will be demonstrated how results from incoherent tunneling measurements and coherent driven dynamics offer complementary approaches to characterizing the noise spectral density [3]. I will also discuss how device design can be used to tune the relative coupling to charge versus flux noise. In particular, by increasing the inductance of our qubits, thus spanning the parametric range from flux qubit to fluxonium, we are able to obtain highly coherent qubits that can be used to probe the electromagnetic environment seen by our quantum processors.

[1] King, A. D., et al. "Coherent quantum annealing in a programmable 2,000 qubit Ising chain." Nat Phys 18.11 (2022): 1324-1328.

[2] King, A. D., et al. "Computational supremacy in quantum simulation." arXiv preprint arXiv:2403.00910 (2024).

[3] Whiticar, A. M., et al. "Probing flux and charge noise with macroscopic resonant tunneling." Physical Review B 107.7 (2023): 075412.

3. Natalie Pearson (PASQAL, France)

Rydberg Blockade Revisited

In Rydberg atom quantum computation, the well known blockade effect is a crucial phenomenon

where an atom in the Rydberg state prevents the excitation of any other nearby atom. Despite its

importance, analysis of this important mechanism is usually simplistic. We extend this analysis to better quantify the blockade phenomenon. Using these insights we explore ways of embedding graphs in Rydberg atom quantum computers by means of register arrangement and pulse sequence, to extend our capabilities for graph-based computation using this hardware.

4. Pol Forn-Díaz (IFAE, spain)

AVaQus project progress

The AVaQus (Annealing-Based Variational Quantum Processors) program is the largest-scale European effort to develop the technology and functionality to operate analog quantum processors as coherent quantum annealers and variational processors. The focus of the project is on both the hardware as well as the theoretical side, to yield a small-scale prototype representing the unit cell of a larger-scale processor that will succeed project AVaQus in the future.

In this talk, I will report on the current state of project AVAQus, focusing on the most recent developments as well as the upcoming challenges. In particular, I will describe the progress developing multi-qubit scalable platforms, the study of theoretical scenarios in search of quantum advantage, as well as the ongoing efforts to produce enabling technology necessary for superconducting qubit technology.

5. Sarah Muschinske (MIT, USA)

Quantum Simulation on a 3x3 Superconducting Qudit Lattice

Quantum simulation with superconducting circuits has largely focused on models of two-level systems such as the hard-core Bose-Hubbard model. However, this choice massively truncates the size of the available Hilbert space – preventing quantum information from being stored in the higher levels of the system and reducing the potential complexity of computation. In this paper, we will discuss the feasibility of performing multi-level analog quantum simulation of the Bose-Hubbard model using superconducting transmon qudits. This approach offers several advantages, including efficient emulation of time evolution (no Trotterization), reduced seepage into the computational subspace (non-Markovian dynamics), and a wider range of Hamiltonians that can be represented. However, with these improvements comes increased complexity in calibration, readout, and tomography protocols. We will give an overview of the design and control of transmon-based Bose-Hubbard emulators. We consider contributions to decoherence and decay of these higher excited states and discuss their impact on many-body behavior. Finally, we will give an outlook for future experiments as we move towards meaningful quantum advantage on noisy intermediate-scale quantum hardware.

6. Naoki Takeuchi (AIST, Japan)

Adiabatic Superconductor Logic: Recent Progress Toward Qubit Controllers

Superconducting digital circuits can operate with low power dissipation at cryogenic temperatures and are thus a promising building block for qubit controllers. Specifically, adiabatic quantum-flux-parametron (AQFP) circuits can operate with an extremely small power dissipation of \sim 10 pW per gate owing to adiabatic switching. In this study, we report recent progress of AQFP logic toward qubit controllers. We present two types of AQFP-based qubit controllers: microwave generator and flux-bias circuit. These circuits can generate qubitcontrol signals (microwave and dc magnetic flux) with a low power dissipation of \sim 100 pW per qubit. Proof-of-concept experiments of these circuits indicate the feasibility of energy-efficient, scalable qubit control using AQFP logic.

This study was supported by the JST FOREST (Grant No. JPMJFR212B) and JSPS KAKENHI (Grants No. JP19H05614 and No. JP19H05615). This study was partly based on results obtained from a project, JPNP16007, commissioned by the New Energy and Industrial Technology Development Organization (NEDO), Japan.

Contributed Talk Abstracts Contributed Talk Abstracts
Jack Raymond : "Quantum error mitigation in quantum annealing " Quantum Error Mitigation (QEM) presents a promising near-term approach to reduce error when

for $\mathbf{F}_{\mathbf{z}}$. $\mathbf{F}_{\mathbf{z}}$ and $\mathbf{F}_{\mathbf{z}}$. We implement $\mathbf{F}_{\mathbf{z}}$. We implement $\mathbf{F}_{\mathbf{z}}$. t. Jack Raymond . Quantum error initigation in quantum anneamig

Quantum Error Mitigation (QEM) presents a promising near-term approach to reduce error when estimating expectation values in quantum computing. Here, we introduce QEM techniques tailored for quantum annealing, using Zero-Noise Extrapolation (ZNE). We implement ZNE through zeronoise quantum annealing, using zero-ivoise Extrapolation (ZIVE). We implement 21VE through zero-
temperature extrapolation as well as energy-time rescaling. We conduct experimental investigations into the quantum critical dynamics of a transverse-field Ising spin chain, demonstrating the successful mitigation of thermal noise through both of these techniques. Moreover, we show that energytime rescaling effectively mitigates control errors in the coherent regime where the effect of thermal anie rescamig enecuvery intigates control errors in the conerent regnific where the effect of their
noise is minimal. Our ZNE results agree with exact calculations of the coherent evolution over a range of annealing times that exceeds the coherent annealing range by almost an order of magnitude.

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2. **Sebastian Leontica** : "Entanglement growth from squeezing on the MPS manifold" \mathbf{r}_{old} and \mathbf{r}_{old} by the Schrödinger equation. A possible solution for the Schrödinger equation for the Schrödinger equation. A possible solution for the Schrödinger equation. A possible solution for the Sch

Finding suitable characterizations of quantum chaos is a major challenge in many-body physics, r manig suitable characterizations or quantum chaos is a major chanenge in many-body physics,
with a central difficulty posed by the linearity of the Schrödinger equation. A possible solution for with a central unficulty posed by the infearity of the Schrodinger equation. A possible solution to
recovering non-linearity is to project the dynamics onto some variational manifold. The classical recovering non-intearity is to project the dynamics onto some variational manriold. The classical
chaos induced via this procedure may be used as a signature of quantum chaos in the full Hilbert space. Here, we demonstrate analytically a previously heuristic connection between the Lyapunov spectrum from projection onto the matrix product state (MPS) manifold and the growth of entanglement. This growth occurs by squeezing a localized distribution on the variational manifold. The process qualitatively resembles the Cardy-Calabrese picture [1], where local perturbations to a movprocess quantativery resembles the Carty-Calabrese picture [1], where local perturbations to a mov-
ing MPS reference are interpreted as bosonic quasi-particles. Taking careful account of the number of distinct channels for these processes recovers the connection to the Lyapunov spectrum. Our results rigorously establish the physical significance of the projected Lyapunov spectrum, suggesting it as an alternative method of characterizing chaos in quantum many- body systems, one that is manifestly connected to classical chaos.

[1] P. Calabrese and J. Cardy, Journal of Physics A: Mathematical and Theoretical 42, 504005 $(2009).$

3. **Roopayan Ghosh** : "n-local catalysts to speed up quantum annealing "
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The quantum speedup obtained to solve optimization problems via quantum annealing is largely bottlenecked by the closing of the energy gap in the changing potential landscape during the quantum anneal. Among several different methods studied to circumvent this hurdle, some of us previously showed th[e](#page-4-0) usage of XX catalyst during annealing increases the gap width¹, that is slows down the gap scaling with system size. In this work we extend the study to non-local catalysts, showing that for different instances of a generalized Maximum Weighted Independent Set (MWIS) problem, an informed choice of n-local (quantum gates involving n-qubits) catalysts can reduce and altogether prevent the closing of gaps during the anneal process. This leverages the symmetry and structure of the problem, but the price to pay is multi qubit quantum gates. The method can be looked upon as a quantum generalization of the multi-spin update process of cluster Monte Carlo algorithms. We discuss the effectiveness and efficacy of the approach under different settings, discussing whether for true quantum advantage one requires scalable multi qubit gates. [1] Feinstein, N., Fry-Bouriaux, L., Bose, S., & Warburton, P. A. (2022). Effects of XX-catalysts

[1] Feinstein, N., Fry-Bouriaux, L., Bose, S., & Warburton, P. A. (2022). Effects of XX-catalysts on quantum annealing spectra with perturbative crossings. *arXiv preprint arXiv:2203.06779*.

and a health annually spectra with pertansance crossings. *arith preprint arith* 2205.00775.
4. **Anna Maria Dziubyna** : "Approximate tensor network contractions for large unitcells quantum annealers \tilde{a}

We develop a heuristic tensor-network-based algorithm aiming to reveal the low-energy spec-trum we develop a heuristic tensor-hetwork-based argorithm anning to revear the low-energy spec- trum
of certain low-dimensional spin-glass systems that encode optimization problems. Inspired by the or certain fow-dimensional spin-glass systems that encode optimization problems. Inspired by the
present-day quantum annealers, we focus on Ising problems defined on quasi-2D lattices with a present-day quantum anneaters, we rocus on ising problems derined on quasi-2D fattices with a
connectivity structure of the current QPU architectures. Our approach combines a branch and bound search strategy with the approximate calculation of marginals via tensor-network con- tractions. Application of the latter numerical techniques to lattices with large unit cells of up to 24 spins, relereppired for the factor numerical techniques to fattices with farge time cens of up to z + spins, referred sparse vant for the current quantum annealing processors, requires a dedicated approach that utilizes sparse structures appearing in tensor network representation, combined with GPU hardware acceleration. We benchmark our approach against D-Wave Advantage QPU and Sim- ulated Bifurcation algorithm for problems defined on Pegasus and Zephyr graphs with up to a few thousand spins. Apart by all the solvers. We observe a good performance of quantum annealers in both these metrics for considered problem classes.

5. **Christopher Baldwin** : "Classical outperforms quantum reverse annealing in fernetic p-spin models) have confirmed that it can indeed avoid phase transitions altogether, thus prolematic first-order phase transitions. The form of the form of the form of the ferromagnetic mean-field models

Adiabatic reverse annealing, first introduced in Ref. [1], is a modification of conventional transverse-field annealing that uses an initial guess at the ground state to attempt to circum- vent problematic first-order phase transitions. Exact solutions for certain mean-field models (the ferromagnetic p-spin models) have confirmed that it can indeed avoid phase transitions altogether, thus providing extremely efficient annealing protocols to locate the desired ground state [2].

In this work [3], we identify the physics underlying the success of (adiabatic) reverse annealing and In this work [5], we identify the physics underlying the success of (adiabatic) reverse annealing and
show that it motivates a corresponding classical reverse annealing algorithm that uses thermal rather inan quantum nuctuations, which we term simulated Teverse annealing (SKA) in contrast to the
above "quantum" reverse annealing (QRA). Thus if QRA is to provide a significant advantage over than quantum fluctuations, which we term "simulated" reverse annealing (SRA) in contrast to the classical algorithms, it must not only succeed in circumventing phase transitions but do so in parameter regimes where SRA fails. Accordingly, we analytically determine the phase diagrams for both we do find regimes in which SRA succeeds while QRA fails). Although discouraging, we also offer QRA and SRA in a wide family of generalized mean-field models $-$ a class sufficiently broad to COMMENTS ON WHY SRA May break down in more realistic problems with rugged energy landscapes, include mean-field approximations to canonical optimization problems such as K-SAT and XOR-SAT — and unfortunately find no regime in which QRA succeeds while SRA fails (interestingly, thus motivating future investigation.

In addition, we complement these results with exact calculations of the real-time dynamics under the addition, we complement these results whilf exact calculations of the real-time dynamics under
both QRA and SRA protocols, confirming that both indeed succeed in reaching the ground state in bout QRA and SRA protocols, communig that bout mueed succeed in reaching the ground state in
O(1) time. This adopts a slightly weaker (but still quite reasonable) definition of "success in O(1) $time$ " — we require that states within arbitrarily small distance of the ground state be reached in O(1) time, but not necessarily the exact ground state itself. Such a definition is more amenable to the tools of statistical mechanics and allows us to perform analytic calculations directly in the thermodynamic limit, which may be of some independent interest.

[1] A. Perdomo-Ortiz, S. E. Venegas-Andraca, and A. Aspuru-Guzik, "A study of heuristic guesses [1] A. Perdomo-Oruz, S. E. Venegas-Andraca, a)

[2] M. Ohkuwa, H. Nishimori, and D. A. Lidar, "Reverse annealing for the fully connected p-spin [2] M. Ulikuwa, H. INISI

[3] C. L. Baldwin, "Classical outperforms quantum reverse annealing in ferromagnetic mean-field models," In preparation.

6. Luca Armando Nutricati : "Using multiple XX-catalysts in quantum annealing to ertifies, which are most problem in the problem in the problem in the problem in \mathbb{R}^n s. Each the mand tradition of structure of the energy landscape, the complexity complement of the complexity co

Whether a quantum annealer is capable of successfully finding the global solution or just a local choice of the embedding anneal embedding is capable of successity in angular grobal solution of just a rocal
minimum strongly depends on the intrinsic characteristics of the problem under study, such as its size, the structure of the energy landscape, the complexity class etc. Along with these specific properties, which are mostly related to the problem itself, other "external" aspects — which rely on how the problem is "prepared" to be solved $-$ could play a crucial role, and include for example the the problem is "prepared" to be solved — could play a crucial role, and include for example the
choice of the embedding, the optimisation of the anneal schedule, the necessity of a reduction up to quadratic couplings etc. Among this variety of aspects, one of the bottlenecks in solving combinatorial optimisation problems using quantum annealers is the emergence of exponentially-closing energy gaps between the ground state and the first excited state during the annealing, which indicates that a first-order phase transition is taking place. As the adiabatic theorem dictates, the total evolution duration of the annealing must scale inversely with the square of the minimum energy gap. For certain problems a first-order phase transition may appear during the anneal, with the minimum energy gap scaling exponentially with the inverse of the system size, ultimately resulting in an exponentially large time required to ensure the adiabatic evolution. For this reason, it is useful to investigate techniques to avoid exponentially closing gaps, or, equivalently, first-order phase transitions. Here we will analyse the effect of using an alternative driver in various maximum weighted independent set (MWIS) problem instances. In Fig. 1 we show the advantage of using the new driver

Figure 1: Avoided first-order phase transitions in the cases of (a) bi-partite and (b) tri-partite MWIS problem
instances. In both cases the II-driver (black) leads to a gap enhancement of two to three orders of magnitude comparison with a usual X-driver (blue) and the use of an XX-catalyst (red).

driver the system does not undergo a first-order phase transition. Finally, an in-depth analysis will be conducted to examine the versatility and potential of the technique across various problem classes, shifting the emphasis away from the MWIS problem as the central theme.

7. **Sudeera Hasaranga Gunathilaka Mastiyage Don** : "Coherent Compressed Sen-
7 sul for the regularised compressed behaviog rial optimisation problems. Most of this interest of this interest has been directed to this interest of the t
Interest of the context of the conte

 T increasingly, quantum-inspired ising machines (QIIMs) have drawn significant attention because rial optimisation problems. Most of this interest has been directed towards implementing QHMs and they have the potential to overcome the difficulties associated with solving large-scale combinatostudying their behaviour, rather than practical applications.

connensional signal of mage using mgmy downscaled measurements. It has attracted significant m-
terest from a wide range of fields and applications. Because L1-norm CS is a convex optimisation To explore the practical applications of OIIMs, we focus on efficiently solving L0-norm compressed sensing (CS) problems with QIIMs. The CS method refers to the reconstruction of a highdimensional signal or image using highly downscaled measurements. It has attracted significant inproblem, many efficient algorithms have been developed that are widely used in real-world applications. The L0-norm CS, however, is believed to perform better than the L1-norm CS. However, optimisation of L0-norm CS is a challenging task since it involves combinatorial optimisation. In the case that L0-norm CS can be solved efficiently by a QIIM, this will pave the way for their practical applications. on CIM (soft- spin-based machine) with an and with an and with an and with an α zeeman term realisation technique. Using the contract of α

In this research, we consider the coherent Ising machine (CIM), one of the most suitable machines for solving optimisation problems in L0-norm CS, since CIMs can construct a densely connected for sorving optimisation problems in E0-norm C5, since Chvis can construct a densery connected
network. We share our proof-of-principle results about how accurately we can perform L0-norm CS network. We share our proor-or-principle results about now accurately we can perform Lo-horm C3
on CIM (soft- spin-based machine) with and without a Zeeman term realisation technique. Using CIMs with chaotic amplitude control (CAC) feedback [3], we demonstrate that they can be used to solve L0-norm CS problems efficiently, outperforming both CIMs without CAC feedback and simulated annealing as well as other practical algorithms such as the Least Absolute Shrinkage and Selection Operator (LASSO) [1]. Furthermore, we show that the performance of CIM with CAC feedback approaches the theoretical limit of L0-norm CS compared to CIMs without CAC feedback [2]. For the real-world data, we have used 64×64 and 128×128 resized MRI images. Results have shown we have used 64×64 and 128×1
ts with CIM with CAC feedback
et al., Sci Rep 13, 16140 (2023) that reconstructio

[1] MDSH Gunathilaka *et al.*, *Sci Rep 13, 16140 (2023)* (2022) [3] Yoshitaka Inui *et al.*, *Commun* [2] Toru Aonishi *et al.*, Quantum Sci. Technol. 7 035013 (2022) [3] Yoshitaka Inui *et al.*, *Commun* Phys 5, 154 (2022)

8. **Natasha Feinstein** : "Escaping local optima in quantum annealing with XX-cou- α . Futusht Fensiem is oscaping focul optima in quantum annealing with αx coa demands on the coherence times of the hardware. A particular problem are the occurrence of pertur-

In quantum annealing, the time needed for a single annealing run to return the optimal solution with high fidelity is dependent on the instantaneous energy spectrum associated with the evolution. Reducing this run-time is of great interest with regards to both the efficiency of the algorithm and the crossings result in exponentially closing gaps in the alliealing spectrum such that exponential an-
neal times are needed to remain in the ground state. One potential route to alleviating this problem demands on the coherence times of the hardware. A particular problem are the occurrence of perturbative crossings which form due to competing local optima of the problem. These perturbative crossings result in exponentially closing gaps in the annealing spectrum such that exponential anis the introduction of an additional catalyst Hamiltonian containing non-stoquastic XX- couplings. In particular, one avenue being explored is the use of targeted catalysts that make use of accessible information about the problem in their construction. In our previous work, we have shown how XXcouplings can be chosen to target a local optimum responsible for a perturbative crossing to alleviate the associated exponential slowdown - either through gap enhancement or by introducing an adbative crossings. Osing scalable graph structures we demonstrate that, by introducing non-stoquas-
tic XX-couplings targeting each of the local optima responsible for a perturbative crossing, we can ditional gap minimum that allows the ground state to be reached through a diabatic path $[1,2]$. Here we show that this approach can be extended to settings where multiple local optima result in perturbative crossings. Using scalable graph structures we demonstrate that, by introducing non-stoquascreate a more efficient path to the ground state. Utilising the fact that polynomial run-times are likely to return precisely the local optima responsible for the perturbative crossings, we discuss the beginnings of a recursive algorithm in which the information needed to select each new coupling could be obtained form beginnings of a recursive algorithm in which the information needed to select each new coupling % could be obtained form the previous annealing run.

[1] N. Feinstein *et al*, arXiv:2203.06779v2

[2] N. Feinstein *et al*, arXiv:2402.13811

9. **Dyon van Vreumingen** : "Gate-based counterdiabatic driving with complexity guarantees" to bise parameters. While the parameter of the such that $\frac{1}{N}$ while one function of the such that $\frac{1}{N}$ $s₁$ time $T₂$, when starting in an initial eigen- state $n₁$ in an initial eigen- state $n₂$, when state $n₁$ and $n₂$ in an initial eigen-

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n of time such that $λ(0) = λ$ _i and $λ(T) = λ$ _f at
n starting in an initial eigen- state $|n(λ$ _i)) of
 $λ$ of H($λ$ f). CD, then, is a method to actively Counterdiabatic driving (CD) has been proposed as an alternative to adiabatic quantum computing (AOC) . In standard AOC , one carries out a time evolution under a gapped time-dependent hamiltonian H(λ (t)), whose parametrisation λ (t) is a function of time such that λ (0) = λ i and λ (T) = λ f at some time T. The result of the driving at time T, when starting in an initial eigen-state $|n(\lambda_i) \rangle$ of H(λ_i), is an approximation of a final eigenstate $\ln(\lambda_f)$ of H(λ_f). CD, then, is a method to actively suppress transitions into other eigenstates in such a pseudo-adiabatic evo- lution. The central object driving through the work of chemical calculations in the work of C^oepaint in the control of the work of CD is the adiabatic gauge potential (AGP) $A(\lambda)$, which is the generator of translations in λ spa $A(\lambda) = -i\partial$. When added to the hamiltonian to construct a λ "counterdiabatic hamiltonian", HCD $(t) = H(\lambda(t)) + \lambda(t)A(\lambda(t))$, transitionless driving is achieved. Finding the AGP, however, is highly nontrivial; hence most research so far has focussed on implementing approximate counterdiabatic driving through bouristics, such as in the work of \tilde{C} engite-ot al. [1]. While these works show in creased fidelity with the final eigenstate as a result of approximate CD, they leave one crucial question unaddressed: what is the added computational complexity of approximate CD as compared to standard AQC, and does the obtained increase in fidelity scale favourably with the added cost? We address this question by making a complexity comparison between a gate-based quantum simulation of standard AQC and a novel gate-based quantum algorithm for CD.

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[2] Sabine Jansen, Mary-Beth Ruskai, and Ruedi Seiler. Bounds for the adiabatic approximation pendent hamiltonians with a trunch that the trunch series. All the distribution series of the series. Application of the series of the ser (2007) .

[3] M[']aria Kieferova['], Artur Scherer, and Dominic W. Berry. Simu- lating the dynamics of time-dependent hamiltonians with a trun- cated dyson series. Phys. Rev. A, 99:042314, Apr 2019.

[4] Guang Hao Low and Nathan Wiebe. Hamiltonian simulation in the interaction picture. 2019.

[5] Pieter W. Claeys, Mohit Pandey, Dries Sels, and Anatoli Polkovnikov. Floquet-engineering counterdiabatic protocols in quantum many-body systems. Phys. Rev. Lett., 123:090602. Aug 2019. [5] Pieter W. Claeys, Mc

[6] Nathan Wiebe, Dominic Berry, Peter Høyer, and Barry C Sanders. Higher order decompositions of ordered operator ex- ponentials. Journal of Physics A: Mathematical and Theoretical, 43(6):065203, jan 2010.

Joseph Cunningham : " Eigenpath traversal by Poisson-distributed phase ran domisation **algorithms** se ran- $\in [0, 1]$, 10 **Ioseph Cunningham** \cdot "Figonpath travereal by Doisson distributed phase ran to. **Boocph cummightin**: Engempum traversur by I obsoin distributed phase fun

and a state in an eigenstate of $H(U)$, we want to prepare the corresponding eigenstate of $H(I)$. Inis is We present a new method for eigenpath traversal: given a path of Hamiltonians H(s), for $s \in [0, 1]$, the same basic goal as Adiabatic Quantum Computing (AQC) and thus it has many applications to the design of quantum algorithms.

Our method is a modification of the Randomisation Method (RM) of $[1]$. In the RM, the eigenstate is tracked by performing "phase randomisation" steps sufficiently often along the path. In order to then the ground state of $H(1)$ can be prepared with arbitrarily small infidelity. One drawback of this perform one step of phase randomisation at some point s, the system is evolved under the Hamiltonian $H(s)$ (as a time-independent Hamiltonian) for some random amount of time. If the right distriburandomisation is performed. This is very different from AQC , where the state evolves smoothly in s aan ri(s) (as a time-muependent riammoman) for some random amount of time. If the right distriction of random times is chosen (this depends on the gap) and this step is performed often enough, method is that the state does not evolve smoothly in s: it jumps at the discrete points where phase

according to the differential equation $d\rho = -iTH(t), \rho$. This equation can then be used to ds derive bounds on the time-complexity, see e.g. [2].

We propose to let the points at which phase randomisation is performed be chosen according to a Poisson process. Then the average state evolves smoothly in s and we show that its evolution is de-

creasing the rate ds termined by $d\rho = \lambda L(\rho)$ for some Lindbladian L. Here λ is the rate of the Poisson process. Thus in-

(i.e. increasing now often the phase randomisation step is performed) improves the nuenty of the re-
nal state, just like increasing T improves the fidelity in AQC. Based on the similarity between the (i.e. increasing how often the phase randomisation step is performed) improves the fidelity of the fi-(i.e. increasing now order the phase randomisation step is performed) improves the ridenty of the final state, just like increasing T improves the fidelity in AQC. Based on the similarity between the two differential eq two differential equations, we can adapt results derived for AQC to our setting.

It is well-known that the performance of AOC depends strongly on the minimal spectral gap Δ m :=

the Hamiltonian H(s). For AQC it is known that the time complexity scales as $O(1/\Delta^{-1})$ if $\frac{R_1}{R_1}$ ds scaling of $O(1/\Delta^{-1})$ using our m version of the RM. In particular, a direct application of our theowhere it and the complexity scales as $O(1/\Delta)$ find the error to the error to the error to $\frac{1}{2}$ find the extension of the error to $\frac{1}{2}$ $t_0 = O(\Delta^{1-p})$ most m holds for all p > 1, see [2]. We show that the same condition reproduces the rem to the Quantum Linear System Problem (QLSP) vields a time-complexity of $O(\kappa \log(1/\epsilon))$. $C(\kappa \log(\kappa \epsilon))$. This time-complexity was previously obtained by an AQC-based algorithm $\left[\frac{4}{7}\right]$. The analysis of such AQC-based algorithms is complicated by the fact that quantum computers cannot where κ is the condition number and ϵ the error tolerance. This is asymptotically optimal and beats the previous state-of-the-art RM-based algorithm [3], whose stated time-complexity is $O(k \log(k/\epsilon))$. This time-complexity was previously obtained by an AOC-based algorithm [4]. The easily simulate a continuously changing Hamiltonian, so some discretisation error is incurred. In order to deal with this, the algorithm in $[4]$ uses the discrete adiabatic theorem, which is arguably more complicated than the usual adiabatic theorem. Our version of the RM side-steps this issue because it is based on the probabilistic execution of a discrete process, which means that there is no discretisation cost.

[2] Sabine Jansen, Mary-Beth Ruskai, and Ruedi Seiler. "Bounds for the adiabatic approximation LIT Sergio Boixo, Emanuel Knill, and Rolando Somma. Elgenpath traversal by phase randomization". In: Quantum Information and Computation 9.9&10 (Sept. 2009), pp. 833–855. doi: 10.26421/QIC9.9-10-7.

[2] Sabine Jansen, Mary-Beth Ruskai, and Ruedi Seiler. "Bounds for the adiabatic approximation with applications to quantum computation". In: Journal of Mathematical Physics 48.10 (Oct. 2007). doi: 10.1063/1.2798382.

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[4] Pedro C.S. Costa et al. "Optimal Scaling Quantum Linear-Systems Solver via Discrete Adiabatic The- orem". In: PRX Quantum 3 (4 Oct. 2022), p. 040303. doi: 10.1103/PRXOuantum.3.040303.

Andrew King : " Computational supremacy in quantum simulation " \mathbf{A} called \mathbf{E}^{\prime} called problem is the such problem is the simulation of non-11. **Anurew King.** Computational supremacy in quantum simulation.

Quantum computers hold the promise of solving certain problems that lie beyond the reach of conventional computers. Establishing this capability, especially for impactful and meaningful problems, remains a central challenge. One such problem is the simulation of nonequilibrium dynamics of a magnetic spin system quenched through a quantum phase transition. State-of-the-art classi-cal simulations demand resources that grow exponentially with system size. Here we show that superconducting quantum annealing processors can rapidly generate samples in close agreement with soconducting quantum annealing processors can rapidly generate samples in crose agreement with so-
lutions of the Schro dinger equation. We demonstrate area-law scaling of entanglement in the model quench in two-, three- and infinite-dimensional spin glasses, supporting the observed stretched-exponential scaling of effort for classical approaches. We assess approximate methods based on tensor networks and neural networks and conclude that no known approach can achieve the same accuracy as the quantum annealer within a reasonable timeframe. Thus quantum annealers can answer questions of practical importance that classical computers cannot.

Ryoji Miyazaki : "Three- and four-body switchable coupler for superconduct-
and **Formish and parity encode the parity encode candidate is the parity encode to part of parity entity** entity e physical-qubit inetwork for this scheme has been proposed.
ing qubits to implement the parity encoding scheme " [talk canceled] $\overline{\text{m}}$ and $\overline{\text{m}}$ are negligible interactions. The interactions in the interactions in the interactions, $\overline{\text{n}}$

for reanzing quantum anneaters. One promising candidate is the parity encountg scheme [1]. The
physical-qubit network for this scheme has been proposed, which is called the parity architecture [2, How we encode interactions of logical qubits in a network of physical qubits is a fundamental issue by. The network is a planar one only with the hearest-heighbor interactions. The interactions, now
ever, are three- and four-body ones. The scheme utilizes such multi-body interactions of physical for realizing quantum annealers. One promising candidate is the parity encoding scheme [1]. The qubits according to the problem instances to be embedded. In other words, one problem requires a 3]. The network is a planar one only with the nearest-neighbor interactions. The interactions, howfour-body interaction of qubits, while another problem requires a three-body interaction for three of the same four qubits. Therefore, in order to implement this theoretical idea, we need a coupler which is not only of multi-body type but can also switch its order of interaction.

We theoretically propose such a coupler for superconducting qubits based on Josephson parametric co use. This is a generalization of a previously proposed coupler for the four-body interaction [4].
The previous circuit is generalized to the one for quarton [5], which is a specific setup of the circuit oscillators. The proposed coupler is a part of four-qubit plaquette and works for three- and fourbody interactions of the qubits. For the three-body interaction, we can select three of the four qubits to use. This is a generalization of a previously proposed coupler for the four-body interaction [4]. constants for the three- and four-body interactions are comparable and can be increased with control called SNAIL [6, 7]. This generalization allows switching from the four-body interaction to the three-body one by modulating the magnetic flux with a certain frequency, which threads the loop consisting of Josephson junctions in the coupler. We numerically demonstrate that the coupling nobs added by our generalization.

 \mathbb{Z} w. Lechner, P. Hauke, and P. Zoller, Sci. Adv. \mathbb{Z} This presentation is based on results obtained from a project, JPNP16007, commissioned by the New Energy and Industrial Technology Development Organization (NEDO).

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Takashi Imoto : "Universal quantum computation using quantum annealing the product of $\frac{1}{2}$ with the transverse-field ising Hamiltonian \tilde{a}

Quantum computation is a promising emerging technology, and by utilizing the principles of quanquantum computation is a promising emerging technology, and by utmzing the principles or quan-
tum mechanics, it is expected to achieve faster computations than classical computers for specific problems. There are two distinct architectures for quantum computation: gate-based quantum comproblems. There are two distinct architectures for quantum computation, gate-based quantum com-
puters and quantum annealing. In gate-based quantum computation, we implement a sequence of quantum gates that manipulate qubits. This approach allows us to perform universal quantum comquantum gates that mampulate quotts. This approach anows us to perform universal quantum com-
putation, yet they pose significant experimental challenges for large-scale integration. On the other hand, with quantum annealing, the solution of the optimization problem can be obtained by prepartional quantum annealing, the solution of the optimization problem can be obtained by prepar-
in the transverse that Concentional menture uses line decisions sith transverse field Line Hamiltoni approximation reliefs transformation of the Hamiltonian of the Hamiltonian, changing from the Hamiltonian, cha
The Hamiltonian of the Hamiltonian of the Hamiltonian, changing from the Hamiltonian, changing the transverse ian, such as those manufactured by D-Wave Inc., achieving around 5000 qubits, are relatively more tan, such as those manufactured by D-wave fitc., achieving around 5000 qubits, are relatively more
amenable to large-scale integration but are limited to specific computations. In this paper, we present a practical method for implementing universal quantum computation within the convenpresent a practical method for imprementing universal quantum computation within the conven-
tional quantum annealing architecture using the transverse-field Ising Hamiltonian. Our innovative and a quantum annealing archiecture using the transverse-rield ising riammoman. Our innovative
approach relies on an adiabatic transformation of the Hamiltonian, changing from transverse fields approach renes on an adiabatic transformation or the rrammoman, changing from transverse riefus
to a ferromagnetic interaction regime, where the ground states become degenerate. Notably, our proposal is compatible with D-Wave devices, opening up possibilities for realizing large-scale gatebased quantum computers. This research bridges the gap between conventional quantum annealing and gate-based quantum computation, offering a promising path toward the development of scalable quantum computing platforms.

14. **Arthur Braida** : "Tight Lieb-Robinson Bound for approximation ratio in quantum annealing is to let a quantum system evolve along a trajectory according to the Ω is the direction of the analog part of the gate-based model (comparable to approximation factor in

Quantum annealing (QA) is one of the two promising frameworks for quantum computing that may end with a quantum advantage because it seems well-suited to solve combinatorial optimization problems. It stands for the analog part of the gate-based model (comparable to QAOA). The goal of quantum annealing is to let a quantum system evolve along a trajectory according to the Schrödinger equation. Comparing to QAOA, there exist few analytical results to attest the theoretical performance of QA. The main one is the adiabatic theorem which states that by allowing a runtime inversely proportional to the square of the minimum spectral gap, it ensures a constant probability of observing the optimal solution. The continuous part of QA makes the equations very difficult to manipulate. cuit to manipulate. $\frac{1}{\sqrt{2}}$ is limited by the instance is limited by the computational power to solve to solve

A recent study from [1] suggests that QAOA, even in the NISQ era may bring a quantum advantage A recent study from [1] suggests that QAOA, even in the ixi5Q era may bring a quantum advantage
over classical algorithm for approximation in optimization problem. Several numerical studies like reful suggests that QA performs wen compared to QAOA however, numerical studies in QA are
rarely convincing as the size of the instance is limited by the classical computational power to solve [2] suggests that QA performs well compared to QAOA however, numerical studies in QA are the Schrodinger equation or by the largest available quantum annealer. To tackle this comparison, cus on the approximation ratio or MaxCut over cubic graph. With one layer, it is formally proven
that QAOA achieves a ratio of 0.6925. In [3], the authors manage to formally prove that constant some researchers tried to develop new mathematical tools to output an analytical value evaluating the performance of QA. As it has been widely used to benchmark metaheuristics, we choose to focus on the approximation ratio of MaxCut over cubic graph. With one layer, it is formally proven sorins QAOA with an approximation ratio or 0.0500. The unriculty in comparing QA to QAOA or
other local algorithm is that QA is non-local by nature. However, we can still make a local analysis time OA achieves a ratio of 0.5933. This last result has been improved by Banks, Brown and Warburton to 0.6003. Above all, Hastings in [4] shows a simple classical local algorithm that outper- I um system. This means that, if the time of the evolution is short enough then two distant sites in the commutativity \mathbb{R}^n forms QAOA with an approximation ratio of 0.6980. The difficulty in comparing QA to QAOA or S ervation can be used to anaryze QA as a focal algorithm to foramily compare approximation ratio of it. In 1972, Lieb and Robinson showed the existence of an information speed limit inside quanquantum system are correlated up to an exponentially small factor in term of the distance. This obvalue of other local algorithms.

In this work, we develop a new tighter Lieb-Robinson (LR) bound by using the commutativity graph from [5] over general regular graphs. This LR bound is adapted to QA applied to MaxCut. 0.7020. This value shows that constant time QA outperforms QAOA with one layer and the best We also slightly modified the standard initial Hamiltonian with a free parameter in front of it. It appears that without changing the algorithm by itself, it helps to improve the tightness of the analysis. Eventually, we end up proving that a 1-local analysis of QA brings the approximation ratio above

n are shows that constant time QA outperform
1-local algorithm.
1-local algorithm.
1.
py, Jonathan Wurtz, Cody Poole, Mark Saf
ncy thresholds for the quantum advantage
npj Quantum Information, 9(1) :73, 2023. $\begin{bmatrix} 2 & 3 & 4 & 5 & 6 \end{bmatrix}$ [1] Danyio Lykov, Jonathan Wurtz, Cody Poole, Mark Saffma *LIJ* Danyio Lykov, Jonathan Wurtz, Cody Poole, Mark Safiman, 10m Noel, and Yuri Alexeev.

Sampling frequency thresholds for the quantum advantage of the quantum approximate optimization al- gorithm. *npj Quantum Informati* Sampling frequency thresholds for the quantum adva

 $\begin{bmatrix} 3 & 7 & 1 & 3 & 1 & 3 & 3 \end{bmatrix}$ [2] Eiljan Peloiske, Andreas Bartschi, and Stephan Eidenbenz. Quantum annealing vs. QAOA :
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**David López-núñez : "Coherent Flux Qubits for Quantum Annealing "

David López-núñez : "Coherent Flux Qubits for Quantum Annealing "** coherent quantum annealers. The Hamiltonian design of the flux qubits is focused on reducing the flux qub

The main work presented in this talk is the initiation of an alternative path towards building coherent quantum annealers based on superconducting flux qubits. An uncoupled flux qubit device is designed to benchmark qubit coherence that will be used as the building block for future iterations of coherent quantum annealers. The Hamiltonian design of the flux qubits is focused on reducing the persistent current to reduce flux noise susceptibility. The processor is designed so that common coherence benchmarking experiments, such as T1 and T2 measurements, are performed. The physical design has been iterated to understand the role the qubit loop configurations and the qubit frequency have on the coherence times.

The flux qubits are measured inside a dilution refrigerator, where special focus is put on magnetic shielding. Spectroscopy measurements provide initial information on the qubit parameters and quality. Coherent control of flux qubits is achieved in the form of coherent Rabi oscillations, which are repeated for many flux operation points to understand the noise mechanisms. Decay times of 40us are shown, which are among the best results for flux qubits. However, the qubit coherence is low, leading to coherence times shorter than 20ns, probably because of flux noise. Moreover, due to device imperfections, the qubits could not be characterized at their optimal coherence conditions, which otherwise corresponds to the initial configuration for quantum annealing. In summary, the coherence results of the flux qubits measured represent an important benchmark in the development of a coherent quantum annealer and provide very valuable information on how to improve the flux qubits for future iterations of quantum annealing processors.

Ana Palacios : "A scalable 2-local architecture for quantum annealing of all-
Fig. 4. μ -allems of interest can be mapped. However, the hard-wiring of a fully connected graph, a general re- α and Tandratic Unconstruction optimization of α and α is which seemed binary optimizations, to which seems the several probability of α .

Quantum annealers are often used as application-specific devices intended to harness quantum ef-Quantum annearers are orien used as apprication-spectric devices intended to narness quantum erquirement for these applications, is not scalable the to crossiant and packing issues within the chip.
We present a scalable architecture to embed an all-to-all connected Ising model within another Ising rects to sorve clas- sical optimisation tasks. This is due to the connection between the ising mode and QUBO (Quadratic Unconstrained Binary Optimisation) formulations, to which several problems of interest can be mapped. However, the hard-wiring of a fully connected graph, a general requirement for these applications, is not scalable due to crosstalk and packing issues within the chip. model defined on a graph of degree $d = 3$ containing exclusively 2-local interactions. This essentially amounts to an efficient braiding of logical chains of qubits which can be derived from a description of the problem in terms of triangles, naturally linking to a family of equivalent formulations of an Ising instance $[1]$.

In addition, we devise strategies to address the chanenges or scalable architectures, such as the
faster shrinking of the gap due to the larger physical Hilbert space and the scaling of penalty suenguis [2, 5], in the spectric context of our architecture. These strategies consist on selecting
driver Hamiltonians that are more suited to the symmetries of the logical solution space, extending In addition, we devise strategies to address the challenges of scalable architectures, such as the I formagnetic X Y Z drivers by examining the scaling of the overlap of the ground state with the logical strengths [2, 3], in the specific context of our architecture. These strategies consist on selecting dard transverse field driver, which relates to a mitigation of the penalty strengths required to find ideas from [4, 5]. In particular, we study the impact of ferromagnetic transverse-field Ising and fersubspace, finding a feasible improvement of around an order of magnitude with respect to the stanthe logical ground state.

Ising Problem for Anti-Crossing Mitigation in Quantum Annealing, Journal of the Physical Society In summary, the nereby prop

tsing Problem for Anu-Crossing Mittgation in Quantum Annealing, Journal of the Physical Society [1] T. Fujii, K. Komuro, Y. Okudaira, and M. Sawada, Eigenvalue-Invariant Transformation of

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 \mathcal{F} anneal in the all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-
Later models, and are all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-to-all-

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Matthias Werner : "Quantum simulation of 1D-fermionic systems with Ising
17. **Matthias Werner** : "Quantum simulation of 1D-fermionic systems with Ising Hamiltonians" ever, many of the set of the set of the discrete on Ising-type Hamiltonians with transverse $\frac{m}{t}$ fields, as the native to are native to quantum hardware platforms like superconducting flux quantum hardware p
The neutral flux quantum hardware platforms like superconduction flux quantum hardware platforms in the neutral

In recent years, programmable, analogue quantum simulators have become capable of simulating neias, as these are hative to quantum hardware platforms like superconducting fiux qubits or neutral quantum critical phenomena in many- body systems, including dynamical phase transitions. However, many of these quantum simulations are focussed on Ising-type Hamiltonians with transverse atoms. The simulation of 1D-systems of spinless fermions, or quantum spin chains, poses a challenge to these platforms due to the lack of non-stoquastic couplings or limited control thereof.

We propose a method to simulate the time- evolution of certain spinless fermionic systems in 1D using simple ising-type i familionians with local transverse fields. Our method is based on domali-
wall encoding [1], which is implemented via strong (anti-)ferromagnetic couplings |J|. We show that simulation of certain fermionic many-body systems accessible to contemporary analogue quantum using simple Ising-type Hamiltonians with local transverse fields. Our method is based on domainin the limit of strong [J], the domain- walls behave like fermions in 1D. This approach makes the hardware that natively implements Ising-type Hamiltonians with transverse fields. The Ising Hamiltonians are 1D chains of spins with nearest-neighbor and, optionally, next-nearest-neighbor interactions.

 U discuss the feasibility of the approach to yield experimental results in the near function \mathcal{U} As a such as phase diagrams and dynamical phase transitions.

Ultimately, we discuss the feasibility of the approach to yield experimental results in the near future. α and α of α 045004 (2019) (20

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18. **Seiya Miyamoto** : "Efficient sampling in a glassy phase using a machineto. Strightly and the carrier sampling in a glassy phase asing a machine learning-assisted Markov Chain Monte Carlo Method " reditinity-assisted markov Gilani mone Gario method

The Markov chain Monte Carlo (MCMC) method can in principle sample from arbitrary probability The Markov chain Monte Carlo (MCMC) method can in principle sample from arbitrary probability
distributions. However, in practice, it often suffers from drastic slowing down in the dynamics and takes a very long time to obtain statisticany independent samples. In physics, such situations
emerge when the system is in the vicinity of a phase transition or in a glassy phase. The details of $\frac{f}{f}$ are ay namical rules in MCMC algorithms strongly control the dynamical slowing down, and, for takes a very long time to obtain statistically independent samples. In physics, such situations some ferromagnetic models, highly efficient non-local algorithms can eliminate the slowing down at the critical points [1,2]. It is thus crucial to devise a better dynamical rule, even if it is 'unphysical', for efficient sampling.

rule. While several non-local algorithms have been proposed for spin glass models [3,4], they are
not efficient enough to access the equilibrium spin glass phase in large systems. Recent studies [5,6] encouraging results. However, in computationally hard, glassy systems with a random first-order transition, they fail to approximate the equilibrium distribution, they fail to approximate the extensive to a
the United Second 200 local electricity is also approximately in even such a contract of the extension of the transition, they fail to approximate the equilibrium distributions, resulting in even slower dynamics have used machine-learning techniques to find non-local Monte Carlo moves, with promising and encouraging results. However, in computationally hard, glassy systems with a random first-order than naive local algorithms $[7,8]$.

Here, we test a machine-learning-assisted algorithm [5–7] on a mean-field model with a full-step mapped onto the maximum independent set of the graph. We find that the algorithm accelerates Here, we test a machine-learning-assisted algorithm [5–7] on a mean-field model with a full-step replica-symmetry breaking transition, another class of glass transitions. We adapt the antiferromagnetic Ising model in a ma replica-symmetry breaking transition, another class of glass transitions. We adapt the antiferromagnetic Ising model in a magnetic field on an Erd" σ s-R envi random graph, with its ground state

global optimum. We will further discuss how the performance of the algorithm depends on the

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Joseph Vovorsh : "Amorphous quantum magnets in a two- dimensional Ryd-
19. **Joseph Vovorsh** : "Amorphous quantum magnets in a two- dimensional Ryd- $\frac{1}{2}$ being atom array $\frac{1}{2}$ crystalline counterpart, there are still many open questions are still many open questions are still many open $\frac{1}{2}$ $\text{coeff}(n)$ is particularly the emergent collective behavior in amorphous materials. This is particular the case in amorphous materials. This is particular that $\text{Diff}(n)$

Amorphous solids, i.e., systems which feature well-defined short-range properties but lack longrange order, constitute an important research topic in condensed matter. While their microscopic structure is known to differ from their crystalline counterpart, there are still many open questions concerning the emergent collective behavior in amorphous materials. This is particularly the case in the quantum regime, where the numerical simulations are extremely challenging. In this talk, we instead propose to explore amorphous quantum magnets with an analog quantum simulator. To this end, we first present an algorithm to generate amorphous quantum magnets, suitable for Rydberg simulators of the Ising model. Subsequently, we use semiclassical approaches to get a preliminary insight of the physics of the model. In particular, we calculate mean-field phase diagrams, and use the linear-spin-wave theory to study localization properties and dynamical structure factors of the excitations. Finally, we outline an experimental proposal based on Rydberg atoms in programmable tweezer arrays, thus opening the road towards the study of amorphous quantum magnets in regimes difficult to simulate classically.

Jonathan Pritchard : "Demonstration of weighted graph optimisation on a
ability to generate large numbers of identical and high quality σ tions the canonical mention can natively implement graph problems in the maximum independent maximum in \mathbf{r} \mathcal{L} . by continuity atoms to highly excited Rydberg states with strong dipole interaction of the m

Neutral atoms have emerged as a powerful and scalable platform for quantum computing, offering the ability to generate large numbers of identical and high quality qubits in reconfigurable arrays. By coupling atoms to highly excited Rydberg states with strong, long-range dipole-dipole interac- α assical opumisation protocols. To extend this approach to explore opumisation of a wider class of tions this system can natively implement graph problems [1] including finding the maximum indeproblems including weighted graphs and quadratic unconstrained binary optimisation (QuBO), it is pendent set (MIS) on a unit disk graph [2], providing a route to performing analogue optimisation of real problems however with large systems required to reach a regime competitive against current necessary to introduce locally addressed light-shifts to enable weighting of nodes during the optimisation process, enabling problems to be encoded with at worst a quadratic resource overheads [3].

In this tark we present work to develop a large-scale system for quantum computing and annealing
and show preliminary results highlighting our ability to implement small-scale demonstrations of In this talk we present work to develop a large-scale system for quantum computing and annealing, weighted graph optimisation using programmable local light-shifts across an atom array [4]. We introduce a hybrid annealing process combining global addressing with ramped light shifts, and outline prospects for scaling this approach to larger graph problems as a potential pathway to quantum utility.

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21. Tomohiro Yamaji : "Parametric dependence of effective local fields of Kerr parametric oscillators zi. **Tomomo Tunaji**. Tunincine dependence of checuve foculated of the KPOs

Kerr parametric oscillators (KPOs), whose Kerr nonlinearities are larger than their photon loss computing [1]. We are developing Josephson parametric oscillators (JPOs) operating in the KPO
regime for application to quantum annealing [2]. The JPOs have two self-oscillating states with a rates, are interesting from the viewpoint of quantum optics and potential applications in quantum *i*n signal to the Viewpoint of quantum optics and potential applications in quantum developing Josephson parametric oscillators (JPOs) operating in the KPO o quantum annealing [2]. The JPOs have two self-oscillating state the oscillation frequency (phase local magnetic phase local magnetic signal corresponds to the local magnetic signal μ relative phase of 0 or p, which can be utilized as a qubit. The oscillation probabilities of the two states can be controlled *in situ* by applying an input signal to the JPOs with the same frequency as the oscillation frequency (phase locking). The phase-locking signal corresponds to the local magnetic fields of an Ising Hamiltonian.

In the classical regime, where the photon loss rate is larger than the Kerr nonlinearity, a JPO can be phase-locked perfectly by applying a weak signal such that the mean photon number in the resphase-locked perfectly by applying a weak signal such that the mean photon number in the res-
onator is about unity [3]. While in the KPO regime, the phase locking of the JPO requires a signal onator is about unity [5]. While in the KIO regime, the phase locking or the 31O requires a signal
stronger than the oscillation output [2]. We study parametric dependence of the phase locking in the KPO regime using an impedance-matched Josephson parametric amplifier [4]. The measured parametric dependence is consistent with the numerical calculation (Fig. 2) and can be understood by considering the steady state of JPO under phase-locking signals. We also discuss the optimization
of the phase locking and quantum annealing using KPOs 2 T. Y. Phys. Rev. Phys. Phys. Rev. A 105, 023519 (2023), T. Y. Yamaji et al., Phys. Rev. Applied 20, 105, 023

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22. **Kunal Prakash Vyas** : "Investigating scaling properties for quantum anneal tems. One such Hamiltonian is the Fermi-Hubbard Hamiltonian. We investigate the scaling complexity for the contraction and the model and gradient contract out using the kinetic energy part as the direction m_{min}

hens. One such riannitomall is the remi-ritiopard riannitomal. We investigate the scaling com-
plexity for the quantum annealing process carried out using the kinetic energy part of the Hubbard plexity for the quantum annealing process carried out using the Kinetic energy part of the Flubbard
model as driving Hamiltonian for ground state calculations. The way we do this is by studying the Quantum annealing can help in finding the ground state of Hamiltonians describing many body systems. One such Hamiltonian is the Fermi-Hubbard Hamiltonian. We investigate the scaling comgaps between the ground state and the 1st relevant excited state that participates in the diabatic evo-
lution for a 1-dimensional system. The behavior of these gaps with increasing system size could rution for a 1-unnensional system. The behavior of these gaps with increasing system size could
hint at polynomial scaling of required annealing time for finding the ground state of the Hubbard gaps between the ground state and the 1st relevant excited state that participates in the diabatic evopreparation for the quantum annealing strategy under study. Information on the complexity coupled Hamiltonian. We also try to extend this idea to a Hubbard model ladder to learn more about the scaling behavior of gaps relevant to adiabatic evolution. Further, we discuss about initial state with an efficient way of preparing the initial state could bolster our hopes for using adiabatic quantum computation for solving correlated many-body Hamiltonians.

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23. **Robert Banks** : "Motivating continuous-time quantum optimisation without zo. **Robert Dams**, motivating commands and quantum optimisation without

Fecourse to the adiabatic theorem
Continuous-time quantum algorithms for combinatorial optimisation problems have previously
been motivated by the adiabatic principle. A number of continuous-time approaches have moved
ward Continuous-time quantum algorithms for combinatorial optimisation problems have previously
been motivated by the adiabatic principle. A number of continuous-time approaches have move
wards exploiting dynamics, no longer ph Continuous-time quantum algorithms for combinatorial optimisation problems have previously commuous-unic quantum argonumis for comonatorial optimisation problems have previously wards exploiting dynamics, no longer physically motivated by the adiabatic principle. In this prewatus exploiting uyuannes, no longer physically mouvated by the adiabatic principle. In this presentation we take rianck's principle as the underlying physical motivation for these approaches
Planck's principle states that the energy of an isolated system cannot decrease as the result of a cyclic process. By applying this principle to multi-stage quantum walks, we show that the approximation ratio can only improve with each stage, motivating monotonic schedules in quantum annealing far from adiabaticity. We also show that the approximation ratio decreases in reverse quantum annealing.

24. **Nick Chancellor** : "A thermodynamic approach to optimization in complex $\frac{1}{2}$ ing protocol in which we repeated with with a contact with a contact with a contact with a contact with a colder auxiliary system of interest in colder auxiliary system of interest in contact with a colder auxiliary system \mathbb{Z}^n , in a complex \mathbb{Z}^n is the energy minimization problem into a thus turned in \mathbb{Z}^n

We start by reviewing our recent published work in [1], which considers the problem of finding the re start by reviewing our recent published work in [1], which considers the problem of finding the
energy minimum of a complex quantum Hamiltonian by employing a non-Markovian bath prepared tem. By tuning the internal parameters of the bath, we show that the optimal cooling is obtained in a regime where the bath exhibits a quantum phase transition in the thermodynamic innit. This result to an addiin a low energy state. The energy minimization problem is thus turned into a thermodynamic cooling protocol in which we repeatedly put the system of interest in contact with a colder auxiliary syshighlights the importance of collective effects in thermodynamic devices.

We present this work as part of larger picture of optimisation far from the adiabatic limit, where we present this work as part or iarger picture or optimisation far from the adiabatic filmt, where
transfering energy from the component of the Hamilto- nian which represents the optimisation problem of interest, as originarly studied in [5]. This picture is extended by coupling to an addi-
tional system. We further present recent unpublished work which examines what happens when a abinal system. We further present recent unpublished work which examines what happens when a
bath is also annealed . Specificially we demonstrate a configuration where, in the adiabatic limit, the baut is also anneated in specificially we definitivate a configuration where, in the adiabatic infint, the system can be returned to it's ground state adiabatically even when initialised. We further argue that problem of interest, as originally studied in [3]. This picture is extended by coupling to an addiarguments of [3]. Based on these arguments it is likely that fine tuning is not needed to im- prove algorithmic performance. Furthermore, we show that all the requirements we have here are realistic even far from the adibatic limit, the system can be made compatible with the energy redistribution to implement in an experimental setting, specifically focusing on a superconducting flux-qubit annealer as an example.

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25. **Gregor Humar** : "Implementing a cooling protocol on a programmable quantum annealer " to. Si spin models with the ground state representing the solution complementation $\frac{3}{2}$. a cooling protocol based on the theoretical protocol based on the theoretical proposal \mathcal{A} on a D-Wave quantum annealer. The theoretical proposal \mathcal{A} on a D-Wave quantum annealer. The theoretical proposal problem

Finding low energy states is an important problem for quantum computers not only for preparing r maing fow energy states is an important problem for quantum computers not omy for preparing
states for quantum simulations [1] but also solving optimization problems [2]. These can be mapped states for quantum simulations $[1]$ but also solving optimization problems $[2]$. These can be mapped
to spin models with the ground state representing the solution [3]. We present an implementation of a cooling protocol based on the theoretical proposal [4] on a D-Wave quantum annealer. The a cooling protocol based on the theoretical proposal [4] on a D-wave quantum annealer. The
process is similar to how cooling is achieved in magnetic refrigeration, where a material in contact process is similar to now cooling is achieved in magnetic refrigeration, where a material in contact
with a heat bath in a strong external magnetic field is first decoupled from the bath before removing the field which cools the material below the original temperature. The protocol consists of coupling qubits representing the problem we want to solve to ancillary qubits representing a spin bath. Starting in the fully polarized state, these bath qubits are subjected to an external field, which is ramped ing in the runy polarized state, these bath qubits are subjected to an external field, which is ramped
down during a cycle of the protocol. Each cycle is ended by a measurement of the system and bath qubits and only the system state is kept for the next cycle, removing the energy that the bath exquotts and only the system state is kept for the hext cycle, removing the energy that the bath ex-
tracted from the system. The protocol is implemented by using annealing offsets to realize different values of fields in the system and in the bath qubits. Reverse annealing is used to set the initial state of the system at the beginning of each cycle. By transferring energy from one subset of qubits to another on the quantum annealing processor, this protocol has the potential to prepare better low energy states than other procedures, which are limited by a finite temperature of the processor. Results of the implementation are presented for the transverse field Ising model system implemented on the D-Wave device and compared to quantum annealing.

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26. **Shuta Kikuchi** : "Evaluation and analysis of hybrid method using quantum zo. Shata Kikachi a Evaluation and analysis of hybrid method using quantum
annealing-machine and non-quantum type Ising machine " anneanng machine and non-quantum type ising machine

Quantum annealing (QA) machines are expected as fast and precision solvers for combinatorial opthe multiple solutions are obtained by non-quantum type ising machine. Next, the problem variables timization problems. To improve the performance of QA machines, a hybrid optimization method combined with a QA machine and a simulated-annealing-based Ising machine (non-quantum type solves the problem-size-reduced sub-ising model by fixing variables. The steps are nerated until the Ising machine) was proposed [1]. The hybrid method is implemented in the following steps. First, solution is stable. In the previous study, the hybrid method outperformed the QA machine alone $[1]$. that consistently have the same value across multiple solutions are fixed. Then, a QA machine

In this study, are resolved that, even formers read that, behydronical method for large-scale problems, that, \sim into the QA machine. The results revealed that, even for large-scale problems, the hybrid method α In this study, we evaluated the performance of the hybrid inethod for farge-scale problems that cannot be input into the QA machine. Simulations were conducted using simulated annealing (SA) implemented by the Markov chain Monte Carlo method as a non-quantum type Ising machine, and a D-Wave Advantage as a QA machine. Sub-Ising models were reduced to the size that can be input indicates higher solution accuracy than the preprocessing SA alone. Furthermore, it was found that the performance of the hybrid method is dependent on the sub-Ising model size.

equation, respectively. The results suggested that the sub-ising model size dependency is attributed Therefore, to investigate the reason for sub-Ising model size dependency, we analyzed through simulations using SA and OA, which were implemented by the master equation and the Schrödinger

[1] Shuta Kikuchi, Nozomu Togawa, and Shu Tanaka, "Hybrid Optimization Method Using Simulated-Annealing-Based Ising Machine and Quantum Annealer," Journal of the Physical Society of Japan, vol. 92, 12 (2023): 124002.

_{Japan,} vol. 92, 12 (2023): 124002.
27. **Asa Hopkins** : "Finding Spin Glass Ground States Using Multi- stage Quan- \mathcal{L} . In the motion of multiple without much without without much value \mathcal{L} and \mathcal{L} t ulti-stage quantum walk, which paired with a heuristic for choosing parameters α

Previous work [1] has shown the advantages of quantum walks compared to other methods for findshows improved scaling compared to single stage walks. Additionally, the scaling exponent appears ing spin glass ground states. In principal, such walks can be implemented on currently existing anto scale as a polynomial in the number of stages, leading to an overall algorithm that has polynomial nealing hardware. In this work, multiple quantum walks are chained together without measurement to produce a multi-stage quantum walk, which paired with a heuristic for choosing parameters

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28. **Lasse Gerblich** : "Advantage of multi-stage quantum walks over QAOA as approximations of quantum annealing 20. Eusse Gerbien. Turrantuge of multi-stage quantum wants over QTOTT a

Methods to find the solution state for optimization problems encoded into ising Hamiltonians are a very active area of current research. We compare the quantum approximate optimization algorithm (QAOA) with multi-stage quantum walks (MSQW). Both are types of variational quantum algorithms, where the control parameters usually need to be optimized classically.

recuveness in sorving such problems is derived from the fact that they approximate the annearing Control theory has been used to show that QAOA corresponds to an optimal protocol for solving optimization problems although they are still outperformed by quantum annealing. Indeed, their efparameters for both protocols by comparing how well they approximate a given annealing schedule.

ready for a superform QAOA with a farge number of stages and an optimal choice of the control parameters for both protocols by comparing now well they approximate a given annealing schedule. We employ a similar argument to show that a multistage quantum walk also corresponds to an optimal protocol for optimization problems. Furthermore, we show analytically that we would expect

ready for a small number of stages. Furthermore, we show that the success probability of finding the We also show numerically for random-spin glass ground state problems, which are a good toy quickly make them impractical as we increase the number of stages, since more parameters need to choose secured model for realistic optimization problems, that we would expect MSQW to outperform QAOA alground state with a QAOA protocol is very sensitive to the choice of control parameters. This will be determined more precisely to ensure the effectiveness of OAOA.

Quantum walks on the other hand do not suffer from this problem. The dependence on the control parameters is a lot smoother. Therefore, choosing "pretty good" heuristic values for them will already provide decent results. Furthermore, energetic considerations can also guide us to choose sensible parameters during the evolution, thus further decreasing the need for powerful classical optimization.

29. Yusuke Sugita : "Parameter-range setting of annealing optimization based on Ising model is model with model α the Ising mathematic 25. **I used cought** Fundacted range search of annealing optimization based of

fising machine is a dedicated system for erriclemly solving the ground-state search problem of the
Ising model [1]. As the operation mechanism of the Ising machine, a typical one is the annealing To realize high-speed and energy-efficient calculations for solving combinatorial optimization probising moder [1]. As the operation mechanism of the ising machine, a typical one is the amealing
method such as simulated annealing (SA) [2] and simulated quantum annealing (SQA) [3]. In the lems, the physics-inspired optimization hardware and algorithms have been actively studied. The Ising machine is a dedicated system for efficiently solving the ground-state search problem of the ficient estimation method of the parameters is nighly desired for the rapid application of the ising method, the annealing parameters control disturbances for variable updates in the optimization process. The appropriate annealing parameters depend on the problems to be solved and thus an efmachine to various real problems.

abinty distribution functions that variables forlow. In the method, first, focal optimal solutions are
obtained by a local search [Figure (a)]. Second, the maximum and minimum values of the annealing Here we investigate a method to estimate the appropriate annealing-parameter range from the probability distribution functions that variables follow. In the method, first, local optimal solutions are Ination, the standard deviation of the probability distribution function is referred to as an indicator
of the movement for each variable. Finally, the main annealing optimization is performed with the parameters are estimated so that the variables move from the local optimum solutions at the maximum value and stay in the local optimum solutions at the minimum value [Figure (b)]. In the estimation, the standard deviation of the probability distribution function is referred to as an indicator determined parameters [Figure (c)]. Due to the second procedure, the appropriate annealing-param-, the appropriations of the the state of the the state of the the state of the state of 2022).
4, 363 (2022). eter range is automatically adjusted according to the model co erricients of the
**220, 4, 363 (2022)
220, 671 (1983).** the presentation, we will show the benchmark results of the n SA method using continuous relaxation for variables [4].

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and **Takuya Okuyama** : "Relaxed Momentum Annealing Theory and Applica-
Takuya Okuyama : "Relaxed Momentum Annealing Theory and Applicaco. **Funity of the capability of the capabilities of convention** computational computational computational strategy in the computational strategy and computational strategy and computational strategy and computational stra $\frac{1}{2}$

The effort to solve NP-hard problems spans various fields, including social network analysis, circuit design, and ma- chine learning, challenging the quest for accurate and efficient solutions [1–3]. The The effort to solve NP-hard problems spans various fields, including social network analysis, circuit design, and ma- chine learning, challenging the quest for accurate and efficient solutions [1–3]. The complexity of the complexity of these problems often surpasses the capabilities of conventional computational strategies. Within this context, the Ising model emerges as a promising approach for combinatorial opti- $\sum_{i=1}^{n}$ algorithms to overcome traditional algorithms $S_1 \subseteq \{-1, 1\}$ interconnected in ought couplings S_1 and \mathbb{R}^2 . We have introduced CMOS and \mathbb{R}^2 as an optimization technique that leverages extensive that leverages e

A has been proven that determining the ground state or ising moders is equivalent to sorving combi-
natorial optimization problems [4]. This method has promoted the use of quantum annealing and Indivital optimization problems [4]. This method has promoted the use of quantum annealing and
other algorithms to overcome traditional algorithmic barriers in nonconvex optimization tasks [5– varier argorithms to overcome traditional argorithmic barriers in nonconvex optimization tasks [5–
9]. We have introduced CMOS Annealing as an optimization technique that leverages extensive It has been proven that determining the ground state of Ising models is equivalent to solving combi-Annealing (MA), enables the simultaneous and independent update of an spins in a fully connected.
In the annual decoders parallel computations in digital circuits [7]. A vital implementation of this technique, Momentum Ising model. MA is capable of approximating the ground state of Ising models with up to 100,000 variables and achieves a speedup of 250 times compared to conventional simulated annealing [10].

The challenge of integrating both binary and continuous variables into optimization methods, such as annealing methods, is substantial. In our presentation, we introduce an annealing approach dees varia
, we i
∈D_i signed to tackie a spectric minimization problem, which includes rinding the ground state or the
Ising model as a creatiol asset $\lim_{n \to \infty}$ for a special case.

min H(s), (2) $s_i \in D_i$ $\min_{\mathbf{r}} \mathbf{r}_{\mathbf{t}}(\mathbf{v}, \mathbf{t})$ of \mathbf{t}

Here, D_i defines the domain for each variable, which could be either $\{-1,1\}$ for binary variables or trated by MA, we incorporate auxinary vectors x and y. The solution to the finital problem is de-
rived from the average of the solutions for x and y. This novel formulation avoids direct interactions $[-1,1]$ for continuous ones. To enable simultaneous spin updates in the fully connected model facilitated by MA, we incorporate auxiliary vectors x and y. The solution to the initial problem is debetween x components, allowing for probabilistic and independent element updates. The same principle is applied to y. Consequently, regardless of the sparsity of J, this new problem structure per-
mits probabilistic updates of all variables in two phases, reflecting the MA method. In this presentation, we will further discuss the practical application outcomes of this algorithm, termed relaxed MA.

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Tatsuhiko Shirai : "Compressed space engineering for constrained combina-
31. **Tatsuhiko Shirai** : "Compressed space engineering for constrained combina studies have demonstrated the effectiveness of tailoring mixer terms for exploring lower-energy sofear **fuguinty simular**. Compressed space engineering for constrained comonia

Combinatorial optimization has attracted much attention as a promising candidate to achieve quancombinatorial optimization has attracted much attention as a promising candidate to achieve quan-
tum speedup. Constraints of a combinatorial optimization problem separate the solution space into tum speedup. Constraints of a comomatorial optimization problem separate the solution space mio
feasible solution space with low energies and infeasible solution space with high energies. Recent studies have demonstrated the effectiveness of tailoring mixer terms for exploring lower-energy sostudies have demonstrated the effectiveness of tanoring mixer terms for exploring fower-energy so-
lutions within the feasible solution space [1]. However, finding appropriate mixer terms is generally rations within the reasible solution space [1]. However, miding appropriate mixer terms is generally
challenging except for a specific class of constraints such as one-hot constraint. To address the issue of restricted applicability, we develop a method to engineer a compressed space. The compressed space expresses the feasible solution space with a smaller number of qubits than the original soluspace expresses the reasible solution space with a smaller humber of qubits than the original solu-
tion space. Since the portion of feasible solutions in the compressed space is large, searching for lower-energy solutions is easier. We develop a scalable variational method to determine the unitary transformation between the compressed space and the original solution space on gate-based quantum computers. Then we propose compressed Quantum Approximate Optimization Algorithm (compressed QAOA) that searches for the (near)-optimal solutions of a combinatorial optimization problem within the engineered compressed space. The figure below shows the simulation result of applying compressed QAOA to quadratic knapsack problems (QKP) with an inequality constraint.
Compressed OAOA outperforms conventional OAOA. This work provides a new framework to solve various types of constrained combinatorial optimization problems.

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32. **Mayumi Nakano** : "Building a Model Learning Method by Reducing the tion problems. When it is a combined a combine with annual optimization by receiving the Number of Training Data for Factorization Machine with Annealing," ivuniber of Franning Data for Factorization machine with Annealing

Ising machines are a promising method for searching for good solutions to combinational optimization problems. When solving a combinational optimization problem with an Ising machine, the objective function and the constraints must be expressed by the Ising model or, equivalently, by the Unconstrained Binary Optimization (QUBO). However, some combinatorial optimization problems, including black-box optimization problems, cannot be expressed in the QUBO form.

Factorization Machine with Annealing (FMA) has been proposed as a black-box optimization ractorization machine with Annealing (FMA) has been proposed as a black-box optimization
method, which uses machine learning and an Ising machine iteratively [1]. The process of FMA is method, which uses machine rearning and an ising machine nerativery [1]. The process of TMA is
divided into two parts: learning and optimization. In the learning process, a machine learning model caned ractorization iviachine (FW) rearns the input-output relationship of an objective function and
expresses the objective function in the QUBO form. In the optimization process, an Ising machine called Factorization Machine (FM) learns the input-output relationship of an objective function and dataset used in the learning process. These two processes are herated, and the best minimum point searches for a minimum point of the QUBO constructed in the learning process. The searched pairs of the minimum point and the value of the objective function for the point are added to the training is returned as a solution to the black-box optimization problem.

When solving constrained black-box optimization problems with FMA, it becomes difficult to generate feasible solutions, which satisfy the constraints. To generate feasible solutions more frequently, we proposed a new model learning method for FM. In our method, the training dataset quently, we proposed a new model learning method for FW. In our method, the training dataset
used in the learning process is always filled with a certain number of the latest training data. That means that, as the FMA iterations proceed, the old training data are removed from the training dataset. In our study, we chose the traveling salesman problem as the target black-box optimization problem. We showed that by selecting the appropriate number of the latest training data, the proposed model learning method can generate feasible solutions more frequently than the conventional FMA.

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Mārtiņš Kālis : "A hybrid quantum-classical approach for inference on re-
33. **Mārtiņš Kālis** : "A hybrid quantum-classical approach for inference on reco. **Thurup runs**. The you dumant classical approach for increme on re stricted Boltzmann machines "

Boltzmann machine is a powerful machine learning model with many real-world applications, for of a D-Wave machine is a powerful machine rearning moder with many real-world applications, for
example by constructing deep belief networks. Statistical inference on a Boltzmann machine can be carried out by sam- pling from its posterior distribution. However, uniform sampling from such a model is not trivial due to an extremely multi-modal distribution. Quantum computers have the thouer is not trivial due to all extremely multi-modal distribution. Quantum computers have the
promise of solving some non-trivial prob- lems in an efficient manner. We explored the application of a D-Wave quantum annealer to generate samples from a restricted Boltzmann ma-chine. The samples are further improved by Markov chains in a hybrid quantum-classical setup. We demonstrated that quantum annealer sam- ples can improve the performance of Gibbs sampling compared to random initialization. The hybrid setup is considerably more efficient than a pure classical sampling. We also investigated the impact of annealing parame- ters (temperature) to improve the quality of samples. By increasing the amount of classical processing (Gibbs updates) the benefit of quantum annealing vanishes, which may be justified by the limited performance of today's quantum computers compared to classical.

Renichiro Haba : "Diverse solutions via quantum annealing leads to the dis-
Renichiro Haba : "Diverse solutions via quantum annealing leads to the discovery or diverse material compositions p_{ref} metaheuristic optimization technique, leveraging p_{ref} unique properties p_{ref}

Finding optimal solutions among a vast number of choices is a critical frontier in disciplines from since D-wave Systems debuted the first commercial device in 2011. A common use of quantum a
nealing is to solve combinatorial optimization problems, aiming to identify states that minimize a materials science to computational science. Recently, quantum annealing (QA) has emerged as a promising metaheuristic optimization technique, leveraging quantum mechanics' unique properties since D-Wave Systems debuted the first commercial device in 2011. A common use of quantum anspecific objective function. Nonetheless, recent research has revealed another benefit of QA: its su-Wide variety of ingredient compinations. Although data science technologies like machine learning perior ability in diverse sampling [1]. Materials exploration is one of the areas where quantum annealing is expected to be applied. In the material industry, in order to develop high performance materials, it is necessary to find the optimal composition from enormous candidates consisting of a offer high-quality predictions of physical properties on uncharted material candidates, finding optimal compositions based on these predictions and certain constraints remains a daunting task in material informatics. α machine α mathematics.

In this study, we introduce a method to identify components with ideal predictive properties using or exact optimization combined with QA machine [2]. Our experimental results mui- cate that the proposed method outperforms random sampling in search efficiency as shown in Fig. 1a. Moreover, black-box optimization combined with OA machine [2]. Our experimental results indi- cate that the we observed that the D-Wave QA machine yielded the highest number of composite solutions as illustrated in Fig. 1b. We found that the diversity of samples produced by OA in each step of blackbox optimization surpasses that of traditional computational meth- ods, such as simulated annealing reation surpasses that or traditive
ptimization and this diversity p
covers new potential uses of Q*li*
diversity is applied practically.
et al. arXiv.2110.10196, 2021. or exact optimization and this diversity positively influences the enumeration of solutions. Our remization and this diversity positively influences the enumerativers new potential uses of QA in materials chemistry and mark versity is applied practically.
al. arXiv.2110.10196, 2021.
et al. Proc. of the 35th Int. Conf

Dario De Santis : "Optimized QUBO formulation methods for quantum com-
Dario De Santis : "Optimized QUBO formulation methods for quantum computing" attention has been paid to the paint of the solve NP-complete problems in the NP-complete problems, which most $\frac{1}{2}$ be formulated as $\frac{1}{2}$ as $\frac{1}{2}$ as $\frac{1}{2}$ and $\frac{1}{2}$ are in-

alate-scale quantum (tyrsty) devices can help to solve optimization problems of any sort. Farticular
attention has been paid to their potential to solve NP- hard and NP-complete problems, which most In recent years, many efforts have been devoted to understand to what extent modern noisy intermediate-scale quantum (NISO) devices can help to solve optimization problems of any sort. Particular stance, an ramous ist is ally s problems [1] can be casted in this form [2]. The possibility to solve
QUBO problems with quantum annealers has been explored widely. Although this technology is not of the times can be formulated as quadratic unconstrained binary optimizations (QUBO). For instance, all famous NP Karp's problems $[1]$ can be casted in this form $[2]$. The possibility to solve expected to solve NP problems efficiently, namely providing exponential speed-ups compared to classical strategies, possible polynomial in-time advantages attract great interest. Indeed, given the importance of the real-world applications involved with the solutions of these hard combinatorial problems, any achievable improvement is well-received. Some examples of these scenarios come from finance, logistics and drug discovery.

logical variables those defining the initial problem. The constrained optimization problem with a
Sider the situation where we aim to solve a generic constrained optimization problem with a
BO solver. The first essent optimization problem with a
dratic unconstrained form. We
ints attached to the original
slack variables. The larger is Consider the situation where we aim to solve a generic constrained optimization problem with a the total number of variables, logical and slack, considered in a Q UBO problem, the narder is to $\frac{1}{2}$ Consider call *logical variables* those defining the initial problem. The constraints attached to the original problem can be enforced in a QUBO form by emploving additional slack variables. The larger is μ_{max} is comparable with the interval variable with the logical variables. In the logical variables. In the set of lo

efficient in several scenarios, their implementation can be extremely inefficient at times. We say
that a method to obtain these formulations is inefficient for a problem whenever it requires too mat a method to obtain these formulations is metricient for a problem whenever it requires too
many slack variables, namely their number is comparable with that of logical variables. In these Whereas the well-established procedures to translate optimization problems as OUBOs [3] can be nany siack variables, namely their number is comparable with that of logical variables. In these
cases, the practical usefulness of QUBO solvers can be highly limited. For instance, if we have a efficient in several scenarios, their implementation can be extremely inefficient at times. We say maximum size for the problems that our solver can receive, any reduction of slack variables allows to increase the total number of logical variables associated to the initial problem, namely larger problems can be tackled. Moreover, the more slack variables we implement, the higher is the connectivity that we require in our QUBO and therefore in the QUBO solver that we aim to use. This consequence can be a very limiting factor for the performances of modern NISO devices.

The main motivation of this work is to unlock the possibility to solve certain classes of optimization The main motivation of this work is to unlock the possibility to sorve certain classes of optimization
problems with quantum annealing. For this purpose, we introduce the iterative quadratic polynomial and master-satellite methods. Whereas the goal of both methods is to provide OUBO forms requiring a minimal employment of slack variables, the former consists in a new paradigm to translate problem constraints into corresponding quadratic penalty forms and the latter allows the simultaneous enforcement of different constraints sharing the same restricted set of binary variables. Noteworthy, non-linear equality and inequality constraints can be treated using the same techniques, without the need to employ slack variables solely to obtain corresponding linearizations. Hence, the same computational effort is required both for linear and non-linear constraints, where no approximate enforcing of the constraints is required.

We follow by approaching an NP-hard problem coming from finance, namely the Max-Profit Balance Settlement prob- lem $[4]$, and we show how to apply our techniques for this optimization problem. In particular, we obtain a drastic reduction of the slack variable necessary to enforce the corresponding inequality constraint. We generate several instances of this problem reflecting some main features of realistic datasets and we show that our methods provide the corresponding QUBO forms by employing around 90% less slack variables, if compared with the QUBO forms obtainable with a standard procedure.

we solve several instances of this optimization problem with two different quantum annealers me
ufactured by D-Wave Systems, Inc., namely the D-wave Advantage system4.1 and We solve several instances of this optimization problem with two different quantum annealers man-Advantage2_prototype1.1, being a prototype of the hext-generation D-wave annealers, has a re-
duced number of qubits but is characterized by higher connectivity and lower noise. Hence, not succed number of quotts but is characterized by inglier connectivity and lower holse. Hence, not
only do we compare the outputs obtained when the QUBO problems are generated with different Advantage2 prototype1.1, which are designed specifically to solve QUBO problems. The annealer Advantage2 prototype1.1, being a prototype of the next-generation D-wave annealers, has a reapproacties, but we also provide a comparison between the performances of these annealers. We
show that, while by using the standard approach the solution quality drops quite fast with the input show that, while by using the standard approach the solution quanty drops quite fast with the input
size, our methods unlock the possibility to consider much larger instances. Indeed, the number of approaches, but we also provide a comparison between the performances of these annealers. We ways significantly more: the multiplicative factor between the successes obtainable with the two successes that we obtain when the QUBO reformulations are generated with our methods are almethods ranges from 7 (smallest instances with Advantage_system4.1) to 184 (largest instances with Advantage2_prototype1.1).

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for \mathbf{A} in the using \mathbf{A} $[2]$ A . Lucas, ising formulations of many NP pro

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Omer Rathore : "Load balancing for high performance computing using
36. **Omer Rathore** : "Load balancing for high performance computing using Here, we present findings from our recent work [1] on the application of quantum annealing to load so. Concernation C. Bodd balancing for their performance computing abing

Here, we present finaings from our recent work $[1]$ on the application of quantum annealing to load With the advent of exascale computing, effective load balancing in massively parallel software applications is critically important for leveraging the full potential of high performance computing systems. Load balancing is the distribution of computational work between available processors. balance two paradigmatic algorithms in high performance computing. Namely, adaptive mesh refinement (AMR) and smoothed particle hydrodynamics (SPH) are chosen as quintessential examples of grid and off-grid based algorithms for computational fluid dynamics.

AMR dynamically adjusts the computational grid to enhance resolution precisely where needed in rear unie, optimizing the anocation or grid resources for complex simulations. SFTT, on the other
hand, eschews the traditional grid framework altogether, instead simulating fluid dynamics through ologies underscore the versathity of algorithms used in fluid dynamics sim- thatfolis, each boasting high perreal time, optimizing the allocation of grid resources for complex simulations. SPH, on the other the interactions of particles that move according to the governing conservation laws. These methodunique advantages. However, due to differences in the underlying data structures, they also present distinct challenges when evenly distributing computational tasks between processors to maximize global efficiency. proval emiclency. The strategic computational challenges more efficiently. We discuss how the strategic im-

tures. This paradigm integrates diverse computing resources, including both classical and quantum Furthermore, our selection of applications is motivated by an appreciation of the evolving high perportance of applications like load balancing becomes evident in this context, as they stand to greatly formance computing landscape, particularly the shift towards heterogeneous computing architecprocessors, to tackle computational challenges more efficiently. We discuss how the strategic imbenefit from such synergistic systems.

While the methodology for obtaining real simulation data to partition is application specific, the proposed balancing protocol itself remains completely general. In a grid based context, quantum annealing is found to outperform classical methods such as the round robin protocol and can potentially remain competitive with more advanced methods such as steepest descent or simulated annealing. For the more complex particle formulation, approached as a multi-objective optimization, quantum annealing solutions are demonstrably Pareto dominant to state of the art classical methods across both objectives. This signals a noteworthy advancement in solution quality which can have a large impact on effective CPU usage.

[1] Omer Rathore, Alastair Basden, Nicholas Chancellor, and Halim Kusumaatmaja. Load balancing for high performance computing using quantum annealing. arXiv preprint arXiv:2403.05278, 2024.

Puya Mirkarimi : "Quantum optimization with linear Ising penalties for cus-
37. **Puya Mirkarimi** : "Quantum optimization with linear Ising penalties for customer data science" Quadratic penalties introduce and equation opening and with this is not penalties for eas $t_{\rm coll}$ and $t_{\rm coll}$ to the performance of a quantum optimizer, especially on near-term optimizer, especially on near-term optimizer, especially on $t_{\rm coll}$

Combinatorial optimization problems that arise in industry are often highly constrained. The standard approach to encoding constraints in quantum optimization is the quadratic penalty method. Quadratic penalties introduce additional couplings and energy scales to a problem's objective function, which can be detrimental to the performance of a quantum optimizer, especially on near-term quantum devices with sparse interaction graphs. We explore an alternative penalty method involving only linear Ising terms, thereby avoiding some of the drawbacks of the quadratic penalty ing results obtained in simulations of quantum annealing and the quantum approximate optimization method, and we apply it to a problem in customer data science. In experiments performed on a D-Wave advantage quantum annealer, we find indications that the linear Ising penalty method can result in improved performance compared to the quadratic method. This is complemented by promisalgorithm with closed-system dynamics. The presented work is based on Refs. [1, 2].

[1] P. Mirkarimi, I. Shukla, D. C. Hoyle, R. Williams, and N. Chancellor, Quantum optimization with linear Ising penalty functions for customer data science, arXiv:2404.05467

[2] P. Mirkarimi, D. C. Hoyle, R. Williams, and N. Chancellor, Experimental demonstration of improved quantum optimization with linear Ising penalties, arXiv:2404.05476

Narendra Hegade : " Digitized Counterdiabatic Quantum Computing " counter diabatic diabatic departum computation, which encodes the chosen in \mathcal{P} bo. I **Narenura riegaue** . Digitized counterdiabatic Quantum computing

a win introduce digitized counterdiabatic quantum computing (DCQC) as a nover paradigm for
compressing digital quantum algorithms. It consists of a suitable digitization of the accelerated I will introduce digitized counterdiabatic quantum computing (DCOC) as a novel paradigm for datic quantum optimization (DCQO) [2, 3]. In particular, I will present an advanced method called
Discussion in the classical methods, or with iterations of process for white iterations in the classical method counterdiabatic dynamics of an adiabatic quantum computation, which encodes the chosen industry use case [1]. I will exemplify DCQC to the class of optimization problems: digitized counterdiabias-field digitized counterdiabatic quantum optimization (bf-DCQO) for tackling combinatorial optimization problems on a digital quantum computer.

any between this combination of the protocols and bias freids offers a way to address farge-
scale combinatorial optimization problems on current quantum computers with limited coherence Along with the selected counterdiabatic (CD) terms in the adiabatic Hamiltonian, we introduce additional bias terms obtained either through classical methods, quantum annealers, or with iterations of DCQO itself. This combination of CD protocols and bias fields offers a way to address largeany classical optimization schemes. In this manner, we overcome the trainability drawbacks faced time. By examining an all-to-all connected general Ising spin-glass problem, we observe a polynomial scaling enhancement in the time to solution compared to both DCQO and finite-time adiabatic quantum optimization. Moreover, the proposed method is purely quantum, eliminating the need for by variational quantum optimization algorithms.

nected spin-glass problem with 33 qubits and a maximum weighted independent set problem with Additionally, bf-DCQO significantly outperforms the quantum approximate optimization algorithm (OAOA) in terms of success probability and approximation ratio. Finally, I will present the experimental results of the proposed method on a trapped-ion quantum computer, tackling a fully conrected spin-giass problem with 33 qubits and a maximum weighted independent set problem with
36 qubits. This represents the realization of the largest quantum computing problem of this nature,
solved on a gate-based quantu 30 qubits. This i

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